# A parameterization of an aerosol physical chemistry model for the $NH_3/H_2SO_4/HNO_3/H_2O$ system at cold temperatures

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**Abstract.** Simple expressions are fitted to the results obtained from ion interaction thermodynamic models for calculating  $HNO_3$  and  $H_2O$  vapor pressures over the  $NH_3/H_2SO_4/HNO_3/H_2O$  system at cold temperatures. The vapor pressure expressions are incorporated into a mass conserving equilibrium solver for computing aerosol compositions in the lower stratosphere and upper troposphere. The compositions calculated from the aerosol physical chemistry model (APCM) are compared against previous parameterizations. The APCM compositions are in better agreement with the compositions obtained from ion interaction models than from other previous formulations of the  $NH_3/H_2SO_4/HNO_3/H_2O$  system. The only advantage of the APCM over the ion interaction approach is that the numerical scheme used in the model is fast and efficient for incorporation into large-scale models. The APCM is used to calculate  $HNO_3$  solubility in ammoniated aerosols as a function of  $HNO_3$ ,  $H_2SO_4$  and  $NH_3$  mass loadings in the lower stratosphere and upper troposphere. While the uptake of  $HNO_3$  by ammoniated aerosols is strongly dependent upon the solution neutrality (or pH), we find that in both the lower stratosphere and upper troposphere a significant fraction of HNO<sub>3</sub> will exist in aerosol solutions near and below the ice frost point irrespective of solution neutrality.

#### 1. Introduction

Thermodynamic electrolyte models are often used for calculating properties of inorganic aerosols in the lower troposphere [Stelson and Seinfeld, 1981; Pilinis and Seinfeld, 1987; Wexler and Seinfeld, 1991; Kim and Seinfeld, 1995; Jacobson et al., 1996; Jacobson, 1999b] and stratosphere [Tabazadeh et al., 1994; Carslaw et al., 1994, 1995b; Weisenstein et al., 1997]. In the lower troposphere, aerosol models are often used in air quality studies to assess the effects of aerosols on health, gasphase partitioning, and visibility. In the stratosphere, aerosol models have been used to simulate the formation and growth of polar stratospheric clouds, which are linked to stratospheric ozone depletion [Solomon, 1999]. However, thermodynamic treatments in large-scale atmospheric models are in general not suited for calculating aerosol compositions in the upper troposphere. Since upper tropospheric aerosols participate in the nu-

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Paper number 2000JD900598. 0148-0227/01/2000JD900598\$09.00 cleation and growth of cirrus clouds and may also be involved in the scavenging of trace gas species, it is important to understand their chemical and physical properties [Kärcher and Solomon, 1999].

Recently, equilibrium aerosol formulations have been incorporated into three-dimensional models to simulate the radiative impacts of aerosols on climate [Adams et al., 1999; Jacobson, 2000]. In this work, an aerosol physical chemistry model (APCM) with an efficient solving scheme suitable for incorporation into large-scale atmospheric models is developed. The APCM is compared against various parameterizations, including results from the aerosol inorganics model (AIM2) of Clegg et al. [1998a].

## 2. Background on Thermodynamic Aerosol Models

#### 2.1. Ion Interaction Approach

The ion interaction approach is originally based on the work of *Pitzer* [1991]. The solution behavior in Pitzer's method is determined by a series of singleion (and water) activity equations that are based on

Table 1. An Coefficients for the Vapor Pressure of HNO3<sup>a</sup>

			1		· · · · · ·			
r	a <sub>0</sub>	$a_1$	$a_2$	$a_3$	$a_4$	$a_5$	a <sub>6</sub>	$a_7$
0.0	12.6923	-0.0325979	-0.071719	0.812331	2.34482	0.000438486	-0.00428908	-0.000709075
0.1	12.4691	-0.0019955	-0.077484	0.727085	2.42858	0.000259723	-0.00442994	-0.000706633
0.2	12.2326	0.0260933	-0.085355	0.644593	2.52519	9.32848e-05	-0.00453571	-0.000694942
0.3	11.9840	0.0513596	-0.094561	0.567113	2.63058	-5.53402e-05	-0.00461276	-0.000678339
0.4	11.7205	0.0725417	-0.104727	0.500403	2.74278	-0.000172282	-0.00466394	-0.000659102
0.5	11.4289	0.0880553	-0.116820	0.450799	2.86783	-0.000241534	-0.00468708	-0.000632214
0.6	11.0837	0.0983196	-0.134271	0.415462	3.02666	-0.000267273	-0.00467490	-0.000579027
0.7	10.6697	0.106923	-0.160517	0.378195	3.24022	-0.000284624	-0.00461980	-0.000480710
0.8	10.2087	0.117295	-0.194708	0.324242	3.50279	-0.000327867	-0.00451884	-0.000342093
0.9	9.73753	0.130202	-0.232997	0.251054	3.79040	-0.000406493	-0.00437409	-0.000184569
1.0	9.28445	0.144365	-0.271608	0.165018	4.07962	-0.000510598	-0.00418932	-2.90206e-05
1.1	8.87335	0.156600	-0.306440	0.080649	4.34517	-0.000611458	-0.00396949	0.00010187
1.2	8.55964	0.157392	-0.327751	0.039034	4.52665	-0.000618013	-0.00372273	0.00015381
1.3	8.33474	0.134427	-0.331971	0.092212	4.60197	-0.000411043	-0.00345255	0.00010694
1.4	7.95881	0.110127	-0.354342	0.142888	4.78919	-0.000208202	-0.00313268	0.00015865
1.5	7.51749	0.0956133	-0.390173	0.143658	5.05815	-0.000120908	-0.00275735	0.00028254
1.6	7.14929	0.0860305	-0.421553	0.118501	5.29796	-0.000104343	-0.00234767	0.000379673
1.7	6.88502	0.0761817	-0.441786	0.092775	5.46818	-0.000107787	-0.00192436	0.00041399
1.8	6.68262	0.0630763	-0.453078	0.080818	5.58436	-0.000101314	-0.00149773	0.00039778
1.9	6.43124	0.0455338	-0.465952	0.086457	5.71507	-7.48198e-05	-0.00106043	0.00038666
2.0	5.72588	0.0242624	-0.524014	0.096528	6.13955	-5.31470e-05	-0.00053951	0.00060982
$\infty^{\mathrm{b}}$	11.5869	0.222271	-0.109389	0.279027	2.83237	-0.00101574	-0.00488557	-0.000704950

<sup>a</sup>  $A_n = a_0 + a_1w_1 + a_2w_2 + a_3\sqrt{w_1} + a_4\sqrt{w_2} + a_5w_1^2 + a_6w_1w_2 + a_7w_2^2$ .

<sup>b</sup> The  $\infty$  represents the ternary system of NH<sub>3</sub>/HNO<sub>3</sub>/H<sub>2</sub>O (i.e., H<sub>2</sub>SO<sub>4</sub> = 0).

thermodynamic properties of mixed solutions. Physical parameters for the fundamental activity relations in an ion interaction model are fitted to laboratory measurements of solute-water mixtures of interest. Thermodynamic electrolyte models based on the ion interaction approach have evolved significantly in recent years by Clegg and coworkers [e.g., Clegg and Brimblecombe, 1990, 1995a, 1995b; Carslaw et al., 1995a; Clegg et al., 1998a, 1998b]. Since the models of Clegg and coworkers, collectively referred to as aerosol inorganics model (AIM2), are fitted to laboratory measurements conducted over a broad temperature range  $(< 200 \text{ to } 328^{\circ}\text{K})$ , they are more accurate than those based on the common binary activity approach described in section 2.2. However, the complex nature of the ion interaction approach makes AIM2 computationally impractical for three-dimensional applications. Here we use AIM2 to generate solution compositions and vapor pressures for the  $NH_3/H_2SO_4/HNO_3/H_2O$ system for a wide range of humidities and temperatures.

#### 2.2. Binary Activity Approach

The second approach first emerged from modeling inorganic aerosols in air quality studies [e.g., Stelson and Seinfeld, 1981; Bassett and Seinfeld, 1983; Saxena et al., 1986]. This method separates out water (via either the Gibbs-Duhem equation or the water equation [Stokes and Robinson, 1966]) and solute activity coefficients. Usually, a mixing rule of either *Bromley* [1973] or Kusik and Meissner [1978] is used to estimate the mixed solute activity coefficients. The equilibrium models based on this approach require only a knowledge of water and solute activity coefficients at a binary level. Thermodynamic models based on the binary activity approach are computationally more efficient but less accurate than the ion interaction approach because the physics of the activity coefficients in the former approach are mainly based on the behavior of binary solutions instead of mixed solution properties used in the latter formulations. To contrast two different thermodynamic treatments, we will update an equilibrium model based on the binary activity approach, EQUI-SOLV II [Jacobson et al., 1996; Jacobson, 1999b], and compare its predictions against results obtained from AIM2. EQUISOLV II applies a well-converged numerical solver scheme to simultaneously solve a large number of equilibrium equations.

#### 2.3. Vapor Pressure Approach

A combination of the two approaches outlined in sections 2.1 and 2.2 is used in this work to develop a fast and accurate parameterization of the  $NH_3/H_2SO_4/$ 

Table	2.	$B_n$	Coefficients	for	$\mathbf{the}$	Vapor	Pressure	of	HNC	) <sub>3</sub> °
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r	b <sub>0</sub>	b1 .	$b_2$	$b_3$	b <sub>4</sub>	$b_5$	b <sub>6</sub>	b <sub>7</sub>
	7679 40		20.2655	250 106	025 20	0.056004	0.07220	0.197376
0.0	-7672.46	67.7280	32.3655	-259.196	-235.30	-0.256884	0.97336	
0.1	-7609.77	50.7068	32.9809	-222.190	-255.43	-0.131368	1.07699	0.201805
0.2	-7541.22	35.7444	34.4801	-189.780	-280.54	-0.026085	1.16451	0.202182
0.3	-7467.24	22.7871	36.6718	-162.151	-309.62	0.059118	1.23812	0.199536
0.4	-7387.11	11.9765	39.4988	-140.155	-342.40	0.122744	1.29896	0.194190
0.5	-7297.89	3.4748	43.2297	-124.533	-380.56	0.163341	1.34688	0.184702
0.6	-7194.29	-3.0516	48.6478	-113.773	-428.87	0.184446	1.38064	0.166777
0.7	-7074.27	-8.6193	56.4190	-103.509	-491.39	0.196255	1.39907	0.136740
0.8	-6944.62	-14.0870	66.1301	-90.235	-565.46	0.207582	1.40222	0.096867
0.9	-6815.18	-19.5849	76.6808	-73.739	-644.27	0.220387	1.39134	0.053245
1.0	-6693.25	-24.6807	87.0206	-56.230	-721.33	0.231211	1.36800	0.011651
1.1	-6584.97	-28.3719	96.0022	-42.304	-789.58	0.230911	1.33398	-0.021582
1.2	-6504.53	-27.8717	100.995	-44.004	-832.71	0.192747	1.29170	-0.031795
1.3	-6449.42	-19.3426	100.908	-77.488	-843.95	0.079582	1.24188	-0.012861
1.4	-6347.58	-8.4904	106.038	-117.723	-886.93	-0.052980	1.17584	-0.022354
1.5	-6203.79	-0.1957	118.040	-143.506	-971.43	-0.156090	1.08786	-0.069346
1.6	-6055.69	4.8272	133.047	-152.440	-1073.17	-0.221191	0.98027	-0.132468
1.7	-5927.18	6.8545	148.101	-146.787	-1173.72	-0.249961	0.85834	-0.195765
1.8	-5829.77	6.2290	161.504	-129.145	-1262.57	-0.245338	0.72779	-0.250386
1.9	-5768.38	3.1411	172.326	-101.381	-1334.05	-0.209351	0.59425	-0.291726
2.0	-5744.66	-2.5250	180.115	-63.965	-1385.49	-0.141562	0.46318	-0.317848
$\infty^{\mathbf{b}}$	-7298.14	-51.7642	35.6176	-55.070	-360.093	0.376735	1.41665	0.252861

<sup>a</sup>  $B_n = b_0 + b_1 w_1 + b_2 w_2 + b_3 \sqrt{w_1} + b_4 \sqrt{w_2} + b_5 w_1^2 + b_6 w_1 w_2 + b_7 w_2^2$ .

<sup>b</sup> The  $\infty$  represents the ternary system of NH<sub>3</sub>/HNO<sub>3</sub>/H<sub>2</sub>O (i.e., H<sub>2</sub>SO<sub>4</sub> = 0).

 $\rm HNO_3/H_2O$  system for atmospheric applications. The vapor pressure approach requires generation of large sets of H<sub>2</sub>O and HNO<sub>3</sub> vapor pressure arrays over the  $\rm NH_3/H_2SO_4/$  HNO<sub>3</sub>/H<sub>2</sub>O system for many assumed compositions and temperatures. The vapor pressure, composition, and temperature fields generated are fitted into simple mathematical expressions similar to those derived by *Luo et al.* [1995] for the H<sub>2</sub>SO<sub>4</sub>/HNO<sub>3</sub>/H<sub>2</sub>O system. The vapor pressure relations are then inserted into a numerical mass-conserving equilibrium solver similar to that used in EQUISOLV II [*Jacobson et al.*, 1996] for gas-aerosol phase partitioning calculations.

#### 3. Model Development

### **3.1.** Aerosol Physical Chemistry Model (APCM)

Luo et al. [1995] have shown that equilibrium partial pressures of HNO<sub>3</sub> and H<sub>2</sub>O over the H<sub>2</sub>SO<sub>4</sub>/HNO<sub>3</sub>/ H<sub>2</sub>O system roughly follow a Clausius-Clapeyron relation of the form  $\ln P = A + B/T$  (where P is pressure, T is temperature, and A and B are constants) for a fixed solution composition. For the quaternary system of NH<sub>3</sub>/H<sub>2</sub>SO<sub>4</sub>/HNO<sub>3</sub>/H<sub>2</sub>O we found a similar behavior for the variation of  $HNO_3$  and  $H_2O$  vapor pressures over the solution.

Assuming that  $NH_3$  and  $H_2SO_4$  reside completely in the condensed phase (i.e., both  $NH_{3(g)}$  and  $H_2SO_{4(g)}$ are negligible), we follow the approach of *Luo et al.* [1995] and define  $w_1$  and  $w_2$  as the weight percents of ammoniated sulfate and nitric acid, respectively, in the solution as follows:

$$w_1 = \text{weight } \% \text{ of } (\text{NH}_4)_r \text{H}_{2-r} \text{SO}_{4(aq)},$$

$$w_2 = \text{weight } \% \text{ of } \text{HNO}_{3(aq)},$$

$$(1)$$

$$r = \frac{NH_3}{H_2SO_4} = \frac{NH_{4(aq)}}{HSO_{4(aq)}^- + SO_{4(aq)}^{2-}},$$
 (2)

where r in  $w_1$  is the fixed mole ratio of ammonia to sulfuric acid, and (aq) is the aqueous-phase species. The r ratio can be considered as the degree of the ammoniated solution neutrality and can take on any values in APCM (including fractions) between 0.0 and 2.0. For example, if r equals 0.0, 1.0, and 2.0,  $w_1$  represents the weight percents of H<sub>2</sub>SO<sub>4</sub>, (NH<sub>4</sub>)HSO<sub>4</sub>, and (NH<sub>4</sub>)<sub>2</sub>SO<sub>4</sub>, respectively, in the solution. Using *Clegg et al.*'s [1998a] model, we first generated a series of arrays of HNO<sub>3</sub> and H<sub>2</sub>O vapor pressures for a wide range of weight percent

Table 3.  $C_w$  Coefficients for the Vapor Pressure of  $H_2O^a$ 

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r	c <sub>0</sub>	c <sub>1</sub>	C <sub>2</sub>	c <sub>3</sub>	C4	C <sub>5</sub>	C <sub>6</sub>	C7
0.0	22.7490	0.0424817	0.0533280	-0.0567432	-0.276555	-0.000621533	-0.000311769	-0.000283120
0.1	22.7669	0.0310031	0.0493720	-0.0270045	-0.265333	-0.000496930	-0.000261912	-0.000257919
0.2	22.7838	0.0223810	0.0457014	-0.0059721	-0.254223	-0.000403349	-0.000224673	-0.000234421
0.3	22.8005	0.0165207	0.0423183	0.0065020	-0.243407	-0.000339520	-0.000198990	-0.000212594
0.4	22.8179	0.0132215	0.0393297	0.0110504	-0.233647	-0.000303073	-0.000184277	-0.000192973
0.5	22.8367	0.0119853	0.0369883	0.0096303	-0.226435	-0.000288721	-0.000180339	-0.000177085
0.6	22.8556	0.0119008	0.0354244	0.0058633	-0.222617	-0.000287316	-0.000186497	-0.000165709
0.7	22.8706	0.0122336	0.0342757	0.0024714	-0.220187	-0.000291401	-0.000201103	-0.000156776
0.8	22.8791	0.0129460	0.0329580	-0.0009337	-0.215893	-0.000300122	-0.000221710	-0.000146949
0.9	22.8818	0.0143893	0.0311831	-0.0063234	-0.208231	-0.000316451	-0.000246179	-0.000134615
1.0	22.8812	0.0168828	0.0289977	-0.0153842	-0.197574	-0.000343217	-0.000273061	-0.000120122
1.1	22.8803	0.0204917	0.0267376	-0.0284821	-0.185921	-0.000381016	-0.000302055	-0.000105483
1.2	22.8846	0.0243030	0.0255578	-0.0417119	-0.179964	-0.000421023	-0.000333949	-9.75631e-05
1.3	22.8946	0.0259818	0.0267000	-0.0449953	-0.186824	-0.000440711	-0.000370519	-0.000103697
1.4	22.8950	0.0265182	0.0275800	-0.0428170	-0.191549	-0.000449334	-0.000409584	-0.000108446
1.5	22.8866	0.0275456	0.0272887	-0.0426774	-0.189011	-0.000462731	-0.000447757	-0.000106656
1.6	22.8752	0.0293967	0.0263037	-0.0461528	-0.182140	-0.000484748	-0.000483688	-0.000101242
1.7	22.8640	0.0319034	0.0251308	-0.0524783	-0.173920	-0.000514332	-0.000517346	-9.51958e-05
1.8	22.8542	0.0347733	0.0240901	-0.0603014	-0.166228	-0.000549028	-0.000549196	-9.03611e-05
1.9	22.8459	0.0376949	0.0233346	-0.0682171	-0.159987	-0.000586063	-0.000579624	-8.75512e-05
2.0	22.8381	0.0403666	0.0228641	-0.0749269	-0.155242	-0.000622692	-0.000608700	-8.66960e-05
$\infty^{\mathrm{b}}$	22.7279	0.0163541	0.0303799	-0.0208336	-0.151376	-0.000354967	-0.000653371	-0.000197299

<sup>a</sup>  $C_w = c_0 + c_1 w_1 + c_2 w_2 + c_3 \sqrt{w_1} + c_4 \sqrt{w_2} + c_5 w_1^2 + c_6 w_1 w_2 + c_7 w_2^2$ .

<sup>b</sup> The  $\infty$  represents the ternary system of NH<sub>3</sub>/HNO<sub>3</sub>/H<sub>2</sub>O (i.e., H<sub>2</sub>SO<sub>4</sub> = 0).

combinations (w<sub>1</sub>, w<sub>2</sub>; where  $0.0\% \le w_1 + w_2 \le 85.0\%$ ) at cold temperatures, ranging from 180 to 270°K. The vapor pressures were then fitted into the following functions:

$$0.0\% \le w_1 + w_2 \le 85.0\%, \quad 180^{\circ}K \le T \le 270^{\circ}K, \quad (3)$$

$$\ln P_{\rm HNO_3} = A_n(r, w_1, w_2) + \frac{B_n(r, w_1, w_2)}{T}, \quad (4)$$

$$\ln P_{H_2O} = C_w(r, w_1, w_2) + \frac{D_w(r, w_1, w_2)}{T}, \quad (5)$$

where  $P_{HNO_3}$  and  $P_{H_{2O}}$  are in mbar. Coefficients  $A_n$ ,  $B_n$ ,  $C_w$ , and  $D_w$  were then parameterized into a mathematical function of  $w_1$  and  $w_2$  similar to that described by *Luo et al.* [1995]. The fitting results are summarized in Tables 1-4 for different r ratios, ranging from 0.0 and 2.0 (r values are separated by an increment of 0.1).

In Plate 1 the calculated vapor pressures using equations (4) and (5) are compared with those obtained from AIM2 for three r ratios (shown in Plates 1a, 1b, and 1c). In general,  $P_{H_{2O}}$  calculated by equation (5) agrees well with the model of *Clegg et al.* [1998a]. However, the agreement of  $P_{HNO_3}$  using equation (4) is rather poor for most compositions. As shown in Plate 1, the calculated  $P_{HNO_3}$  becomes increasingly inaccurate as the solution neutrality (i.e., r) increases. For some cases the difference between estimated  $P_{HNO_3}$  values and those predicted by AIM2 reached as high as a factor of 5. To ensure differences are minimized, Diff terms accounting for the differences (deviations) of the "calculated" values (estimated from equations (4) and (5)) and the "real" values (obtained from AIM2) were added to equations (4) and (5):

$$lnP_{HNO_{3}} = A_{n}(r, w_{1}, w_{2}) + \frac{B_{n}(r, w_{1}, w_{2})}{T} + Diff_{n}(r, w_{1}, w_{2}, T)$$
(6)

$$\ln P_{H_{2}O} = C_{w}(r, w_{1}, w_{2}) + \frac{D_{w}(r, w_{1}, w_{2})}{T} + Diff_{w}(r, w_{1}, w_{2}, T).$$
(7)

For a fixed r,  $w_1$ , and  $w_2$  the Diff correction terms approximately follow simple polynomial functions in temperature:

$$\operatorname{Diff}_{n}(\mathbf{r}, \mathbf{w}_{1}, \mathbf{w}_{2}, \mathbf{T}) = \sigma_{n0} + \sigma_{n1}\mathbf{T} + \sigma_{n2}\mathbf{T}^{2} \qquad (8)$$

r	d <sub>0</sub>	$d_1$	$d_2$	d <sub>3</sub>	d4	$d_5$	$d_6$	d <sub>7</sub>
0.0	-5850.24	21.9744	1.48745	-44.5210	59.6400	-0.384362	-0.644671	-0.208556
0.1	-5849.94	19.8594	1.73322	-39.7129	58.0581	-0.336545	-0.606959	-0.210361
0.2	-5849.15	17.5776	1.92716	-34.3704	56.5456	-0.289409	-0.568784	-0.211815
0.3	-5848.04	15.1511	2.06677	-28.5548	55.1548	-0.243024	-0.530295	-0.212915
0.4	-5846.79	12.6262	2.12829	-22.4377	54.0535	-0.197733	-0.491539	-0.213544
0.5	-5845.48	10.1120	2.06175	-16.4657	53.5369	-0.154503	-0.452501	-0.213397
0.6	-5843.72	7.79449	1.85460	-11.3889	53.6970	-0.115030	-0.413300	-0.212396
0.7	-5840.63	5.79437	1.60810	-7.6468	53.9687	-0.0803938	-0.374279	-0.211123
0.8	-5835.71	4.05810	1.46205	-4.9098	53.5707	-0.0500282	-0.335935	-0.210378
0.9	-5829.33	2.44674	1.47151	-2.5108	52.2170	-0.0225318	-0.298645	-0.210471
1.0	-5822.28	0.84270	1.60302	0.0831	50.1179	0.0033213	-0.262577	-0.211194
1.1	-5815.44	-0.79036	1.74377	2.9989	47.9281	0.0280087	-0.227587	-0.211883
1.2	-5810.31	-2.20859	1.55496	5.1396	47.5949	0.0492844	-0.193207	-0.210535
1.3	-5807.19	-2.71615	0.62592	3.4557	51.4599	0.0605252	-0.158668	-0.204729
1.4	-5801.66	-2.43122	-0.39261	-1.5150	55.7388	0.0628630	-0.124306	-0.198357
1.5	-5792.68	-1.85557	-1.06065	-7.4818	57.9126	0.0612095	-0.091134	-0.193964
1.6	-5781.23	-1.28487	-1.31402	-13.0971	57.6578	0.0586697	-0.059876	-0.191880
1.7	-5768.33	-0.88563	-1.18923	-17.6257	55.2142	0.0571507	-0.030919	-0.191856
1.8	-5754.83	-0.75269	-0.74364	-20.6765	50.9331	0.0578653	-0.004421	-0.193536
1.9	-5741.30	-0.93290	-0.03072	-22.0791	45.1436	0.0615518	0.019549	-0.196612
2.0	-5728.09	-1.44363	0.91523	-21.7913	38.0640	0.0686466	0.040959	-0.200889
$\infty^{\mathrm{b}}$	-5736.70	-0.27462	2.02241	0.5073	20.6699	0.0337051	-0.045872	-0.179320

Table 4.  $D_w$  Coefficients for the Vapor Pressure of  $H_2O^a$ 

<sup>a</sup>  $D_w = d_0 + d_1w_1 + d_2w_2 + d_3\sqrt{w_1} + d_4\sqrt{w_2} + d_5w_1^2 + d_6w_1w_2 + d_7w_2^2$ .

<sup>b</sup> The  $\infty$  represents the ternary system of NH<sub>3</sub>/HNO<sub>3</sub>/H<sub>2</sub>O (i.e., H<sub>2</sub>SO<sub>4</sub> = 0).

$$Diff_{w}(r, w_{1}, w_{2}, T) = \sigma_{w0} + \sigma_{w1}T + \sigma_{w2}T^{2}, \quad (9)$$

where  $\sigma$  are second order polynomial coefficients fitted for a particular combination of  $w_1$ ,  $w_2$ , and r. The effects of including Diff terms in improving the vapor pressure fits are shown in Plates 1d-1f where HNO<sub>3</sub> and  $H_2O$  vapor pressures calculated from equations (6) and (7) (with polynomial coefficients of equations (8) and (9)) are compared against AIM2 results for the same r ratios examined above. As shown in Plates 1d-1f, Diff terms force the calculated vapor pressures of both  $H_2O$ and HNO<sub>3</sub> to agree with AIM2 results. For r values other than 0.0, 1.0, and 2.0, including Diff terms in vapor pressure relations produces a nearly exact agreement with AIM2 results. Thus, for APCM we tabulated polynomial coefficients of  $\sigma$  (lookup tables available as electronic supporting material<sup>1</sup>) for all possible combinations of  $w_1$ ,  $w_2$ , and r. The weight percents of  $w_1$  and  $w_2$  in the lookup tables cover and span the composition spectrum from 0.0% to 85.0% with 1% increments in composition.

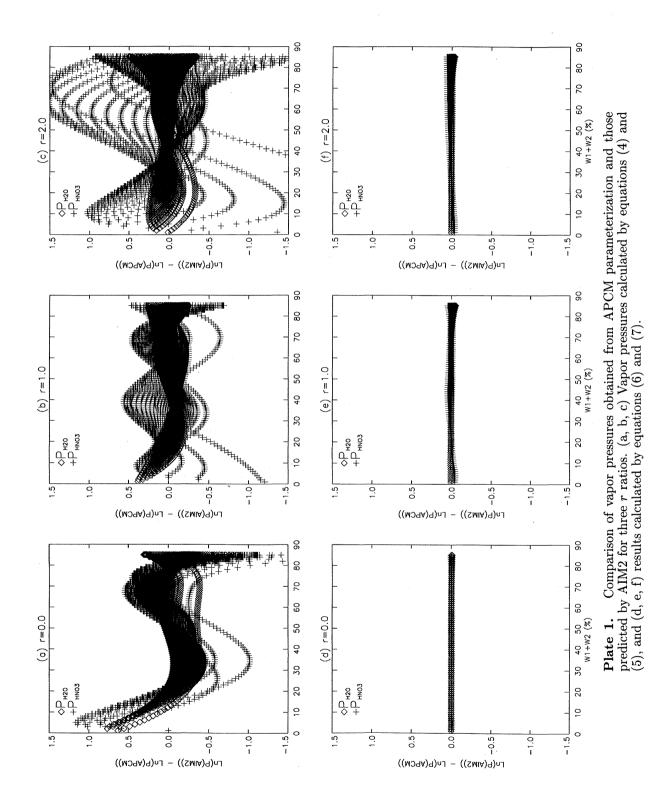
As temperatures cool, solubilities of trace gases, such as nitric acid, increase significantly, thereby depleting gas phase concentrations. Thus it is essential to calculate the distribution of  $HNO_3$  between the gas and aerosol phases. To simulate the gas-aerosol partitioning of  $HNO_3$ , equations (6) and (7) are coupled with mass conservation of  $HNO_3$ :

total HNO<sub>3</sub> = 
$$P_{HNO_3}$$
 + HNO<sub>3(aq)</sub> (in mol/m<sup>3</sup>). (10)

Three unknowns,  $(NH_4)_rH_{2-r}SO_{4(aq)}$  (or  $w_1$ ),  $HNO_{3(aq)}$  (or  $w_2$ ), and  $P_{HNO_3}$ , can be uniquely determined by equations (6), (7), and (10). For this work, the same numerical scheme as that utilized in EQUISOLV II [*Jacobson et al.*, 1996] is applied to solve equations (6)-(10) iteratively.

At iteration steps where  $w_1$  and  $w_2$  are not tabulated in the lookup tables, the bilinear (area weighted) averaging [*Jacobson*, 1999a] is adopted for interpolation. For example, to estimate properties at  $[w_1, w_2] =$ [25.3, 10.8], results obtained from four adjacent points,

<sup>&</sup>lt;sup>1</sup>Supporting lookup tables are available via Web browser or via Anonymous FTP from ftp://kosmos.agu.org, directory "apend" (Username = "anonymous", Password = "guest"); subdirectories in the ftp site are arranged by paper number. Information on searching and submitting electronic supplements is found at http://www.agu.org/pubs/esupp\_about.html.



	А	В	С	D				
	$(NH_4)_2SO_4$							
$0.00 < a_w < 0.30$								
<b>y</b> 1	$5.1229560233e{+}01$	-9.6193408336e-01	1.4720578701e+02	-1.7229112367e-02				
$y_2$	$4.5462384945e{+}01$	$-9.4872028917e{-}01$	$3.3412155166e{+}02$	-2.1124614825e+02				
$0.30 \leq a_w < 0.60$								
<b>y</b> 1	9.7348916855e - 01	-2.7444251275e+00	-1.8966528614e+01	$1.9212056647e{+}01$				
$y_2$	1.0148647694e+00	-2.3916961554e+00	-2.2987950986e+01	$2.1920018865e{+}01$				
$0.60 \leq a_w \leq 1.00$								
<b>y</b> 1	-5.3756389490e+02	1.3556456453e-02	-1.8951248281e+01	5.5692105238e+02				
<b>y</b> <sub>2</sub>	-4.2921431877e+02	1.2875856340e-02	-2.1275933520e+01	4.5075478762e + 02				
		$NH_4$	NO <sub>3</sub>					
$0.00 < a_w < 0.40$								
$\mathbf{y}_1$	1.3165236133e+02	-1.0103415600e+00	-3.1726829968e+01	-1.4042075761e+02				
<b>y</b> <sub>2</sub>	8.1149051028e+01	-9.6213385243e-01	$1.0536766694e{+}02$	-1.5360919227e+02				
$0.40 \leq \mathrm{a_w} < 0.85$								
<b>y</b> 1	8.5782004669e + 01	-1.2795635979e+00	2.9982661442e+01	-1.0776290885e+02				
У2	3.2723371194e+01	-1.3299217851e+00	-8.1245973075e+00	-2.4520261799e+01				
$0.85 \leq a_w \leq 1.00$								
y1	-1.4969599516e-01	-2.0271382540e+01	-2.1065828892e+02	2.0679076661e+02				
<b>y</b> 2	1.6636062187e + 03	$-1.3465428608e\!-01$	$1.8592287444e{+}02$	-1.8496676635e+03				
		$NH_4H_2$	ISO4					
$0.00 < a_w < 0.25$								
<b>y</b> 1	$8.9148556445e{+}01$	-9.2225449765e-01	2.1668157712e + 02	-2.7552338531e+02				
<b>y</b> 2	5.9840874514e+01	-9.3729948048e-01	2.7226224522e + 02	-2.1721584622e+02				
$0.25 \leq \mathrm{a_w} < 0.80$								
y1	$2.8756459725e{+}00$	-2.5507156109e+00	-3.3316860933e-01	-1.8494214012e-03				
$y_2$	$3.5550024880e{+00}$	-2.0910350500e+00	-1.0492539829e+01	7.8923013513e+00				
$0.80 \leq a_w \leq 1.00$								
<b>y</b> 1	-1.4227339985e+00	2.9277231432e+01	-1.7215058088e+01	1.8441662636e+01				
<b>y</b> <sub>2</sub>	$-7.0483234774\mathrm{e}{-01}$	$2.0832745270e{+}01$	-2.3289593929e+01	$2.3901137559e{+}01$				

Table 5. Composition Functions<sup>a</sup>

<sup>a</sup>  $y = Aa_w^B + Ca_w + D$ . Read 5.1229560233e+01 as 5.1229560233×10<sup>1</sup>. All parameterizations are valid for 190° K  $\leq T \leq 260^\circ$  K only. Molality m is calculated by  $m(a_w, T) = y_1(a_w) + (T - 190)$   $[y_2(a_w) - y_1(a_w)]/70$ , where  $a_w$  is water activity (relative humidity expressed in fraction). Composition functions for H<sub>2</sub>SO<sub>4</sub> and HNO<sub>3</sub> in the same mathematical form are given by *Tabazadeh et al.* [1997a, 1997b].

[25.0, 10.0], [25.0, 11.0], [26.0, 10.0], [26.0, 11.0], are weighted and averaged, depending on how far the desired point is located away from the four fixed points. Similarly, a simple linear interpolation (of averaging predictions from two adjacent r ratios) is implemented if initial r falls inbetween the 0.1 increments (see Table 1). Since the intervals are small (1% for  $w_1$  and  $w_2$  and 0.1 for r), the two interpolation methods yield reasonable results that are in close agreement with AIM2 in predicting HNO<sub>3</sub> and H<sub>2</sub>O vapor pressures over the NH<sub>3</sub>/H<sub>2</sub>SO<sub>4</sub>/HNO<sub>3</sub>/H<sub>2</sub>O system.

For conditions where  $H_2SO_4$  mixing ratio is zero we define  $w_1$  and  $w_2$  as the weight percents of  $NH_4NO_{3(aq)}$  and  $HNO_{3(aq)}$ , respectively, and repeat the above steps to obtain vapor pressures and solution compositions for the  $NH_3/HNO_3/H_2O$  system.

#### 3.2. An Update of EQUISOLV II

In addition to APCM parameterization an equilibrium model based on the binary activity approach [Stelson and Seinfeld, 1981; Pilinis and Seinfeld, 1987; Wexler and Seinfeld, 1991; Kim and Seinfeld, 1995; Jacobson et al., 1996; Jacobson, 1999b], EQUISOLV II [Jacobson et al., 1996; Jacobson, 1999b], is updated for application at colder temperatures. The updated EQ-UISOLV II will be compared against AIM2 and APCM. The important equilibrium equations to solve in EQUI-SOLV II for the  $NH_3/H_2SO_4/HNO_3/H_2O$  system are

$$\mathrm{HSO}_{4(\mathrm{aq})}^{-} \stackrel{\mathrm{K}_{\mathrm{HSO}^{-}}}{\longleftrightarrow} \mathrm{H}_{(\mathrm{aq})}^{+} + \mathrm{SO}_{4(\mathrm{aq})}^{2-}, \qquad (11)$$

$$HNO_{3(g)} \stackrel{K_{HNO_3}}{\longleftrightarrow} H^+_{(aq)} + NO^-_{3(aq)}, \qquad (12)$$

	$\alpha_k(T) = \beta_0 + \beta_1 T + \beta_2 T^2$							
Coefficients	$\beta_0$	$eta_1$	$\beta_2$					
		$HNO_3(n=5)$	· · · · · · · · · · · · · · · · · · ·					
$lpha_0$	$1.1550607531e{+}01$	-8.6598065340e-02	1.6100335132e-04					
$lpha_1$	-2.6045646534e+01	1.8570118152e-01	-3.4032931641e-04					
$lpha_2$	$9.6192566768e{+}00$	-6.4314120728e-02	1.1363959743e-04					
$lpha_3$	-1.4330104905e+00	9.3470490393e-03	-1.6321073928e-05					
$lpha_4$	9.5181767818e-02	-6.1140131174e-04	1.0605856919e-06					
$lpha_5$	-2.3568483179e-03	1.4936021604e-05	$-2.5755879088e{-}08$					
		$(NH_4)_2SO_4(n=3)$						
$lpha_0$	$3.4428316946e{+}00$	-4.5328321654e-02	1.0202137529e-04					
$lpha_1$	-1.0030630804e+00	6.0339372495e-03	-1.2162329645e-05					
$lpha_2$	$4.2925175339\mathrm{e}{-02}$	-3.0262230896e-04	5.9904176768e-07					
$lpha_3$	-6.6813265363e-04	5.0031688798e-06	-9.7460258587e-09					
		$NH_4NO_3(n=4)$						
$lpha_0$	2.7037375056e+00	-1.7880171498e-02	2.8497778393e-05					
$lpha_1$	-4.9424685844e+00	2.1621849334e-02	-2.2789057453e-05					
$lpha_2$	9.0584578176e - 01	-4.2947161383e-03	4.7971905516e-06					
$lpha_3$	-7.7554691781e-02	3.7618465057e-04	-4.3345626096e-07					
$lpha_4$	2.4783745998e-03	$-1.2123165585e{-}05$	1.4181424491e-08					
		$NH_4HSO_4(n=5)$						
$lpha_{0}$	9.9601477451e-01	-1.1296951291e-02	1.7452716744e-05					
$\alpha_1$	1.4845831451e+00	-9.3756384679e-03	$1.8989492030e{-}05$					
$lpha_2$	-2.9273249432e-01	1.3469653174e-03	-2.2718917359e-06					
$lpha_3$	$2.1988585949e{-}02$	-8.1477260064e-05	1.0470356871e-07					
$lpha_4$	-7.5397112432e-04	2.2096863963e-06	-1.6051821078e-09					
$lpha_5$	9.7902066682e-06	-2.1560396221e-08	$-3.3248336670e{-12}$					

Table 6. Mean Binary Activity Coefficients<sup>a</sup>

<sup>a</sup>  $\ln\gamma = \sum_{k=0}^{n} \alpha_k(T)m^{k/2}$ , Read 1.1550607531e+01 as 1.1550607531×10<sup>1</sup>. All parameterizations are valid for 190° K  $\leq T \leq 260^{\circ}$  K only. Ion activity for the  $H_2SO_4/H_2O$  system was not parameterized but instead tabulated in the computer codes.

$$\mathrm{NH}_{3(g)} + \mathrm{H}^{+}_{(\mathrm{aq})} \stackrel{\mathrm{K}_{\mathrm{NH}_{3}}}{\Longleftrightarrow} \mathrm{NH}^{+}_{4(\mathrm{aq})}, \qquad (13)$$

where  $K_i$  is the equilibrium constant and (g) and (aq) refer to gas- and aqueous-phase species, respectively. The first dissociation step of  $H_2SO_4$  ( $H_2SO_{4(1)} \rightarrow HSO_{4(aq)}^- + H_{(aq)}^+$ ) is assumed to be complete.

Temperature-dependent water and solute activity coefficients in EQUISOLV II are modified at cold temperatures (190°K-260°K) by parameterizing data from the model of *Clegg et al.* [1998a]. The temperaturedependent water activities of five electrolytes involved in equations (11)-(13), (H<sub>2</sub>SO<sub>4</sub>, HNO<sub>3</sub>, (NH<sub>4</sub>)<sub>2</sub>SO<sub>4</sub>, NH<sub>4</sub>NO<sub>3</sub>, NH<sub>4</sub>HSO<sub>4</sub>) are improved in EQUISOLV II. For consistency with previous work by *Tabazadeh et al.* [1997a, 1997b] the binary solution compositions are parameterized into

m = y<sub>1</sub>(a<sub>w</sub>) + 
$$\frac{(T - 190)[y_2(a_w) - y_1(a_w)]}{70}$$
, (14)

where m is the molality of the binary electrolytes,  $a_w$  is the water activity (relative humidity expressed as fraction), and y is a function of  $a_w$  only (see Table 5). Mean binary activity coefficients  $(\gamma_{ij}^o)$  were parameterized by converting mole fraction activity coefficients (f) of *Clegg et al.*'s [1998a] model into molality base  $(\gamma)$  values according to equation (15), and using  $\gamma_{ij}^o =$  $(\gamma_i^{\nu_+} \gamma_j^{\nu_-})^{\frac{1}{\nu_++\nu_-}}$  [*Pitzer*, 1991]:

$$f_{i} = \gamma_{i} \left( 1 + \frac{M_{w}}{1000} \sum_{k} m_{k} \right), \qquad (15)$$

where  $\nu_+$  and  $\nu_-$  are the stoichiometric coefficients of the binary electrolytes (for example,  $\nu_+ = 2$  and  $\nu_- = 1$ for  $(NH_4)_2SO_4$ ),  $M_w$  is the molecular weight of water, and summation is over all solute species. The mean activity coefficients were then fitted into simple polynomial functions (given in Table 6) of temperature and molality (where  $\alpha_k$  and  $\beta_k$  are polynomial coefficients):

Cases	$r=\frac{NH_3}{H_2SO_4}$	$\mathrm{NH}_3,\mathrm{ppt}$	$H_2SO_4$ , ppt	HNO <sub>3</sub> , ppt	H <sub>2</sub> O, ppm		pm	Illustration		
Stratospheric Simulation ( $P_{total} = 50 \text{ mbar}$ )										
Background	r = 0.0	0.0	500.0	10,000.0		5.0		Figure 1a and 1b		
Volcanic	r = 0.0	0.0	20,000.0	10,000.0		5.0		Figure 2a and 2b		
i.		Upper Trop	ospheric Simula	tion $(P_{total} = 200)$	) mbar )	) 				
Background	r = 0.0	0.0	100.0	100.0	5.0,	50.0,	500.0	Figure 3a		
	r = 1.0	100.0	100.0	100.0	5.0,	50.0,	500.0	Figure 3b		
	r = 2.0	200.0	100.0	100.0	5.0,	50.0,	500.0	Figure 3c		
Polluted	r = 0.0	0.0	100.0	2,000.0	5.0,	50.0,	500.0	Figure 4a		
	r = 1.0	100.0	100.0	2,000.0	5.0,	50.0,	500.0	Figure 4b		
	r = 2.0	200.0	100.0	2,000.0	5.0,	50.0,	500.0	Figure 4c		
		HI	$VO_3$ Uptake ( $P_{to}$	otal=200 mbar)						
Background	r = 0.0 - 2.0	0.0 - 200.0	100.0	100.0	5.0,	50.0,	500.0	Figure 5		
	$\mathbf{r} = \infty$	100.0	0.0	200.0	5.0,	50.0,	500.0			

 Table 7. Model Conditions<sup>a</sup>

<sup>a</sup>  $P_{total}$  is total atmospheric pressure, and the mixing ratios used are atmospheric observations [*Tabazadeh et al.*, 1998; *Laaksonen et al.*, 1997].

$$\ln \gamma_{ij}^{o} = \sum_{k=0}^{n} \alpha_{k}(T) m^{k/2}, \qquad (16a)$$

$$\alpha_{\mathbf{k}}(\mathbf{T}) = \beta_0 + \beta_1 \mathbf{T} + \beta_2 \mathbf{T}^2.$$
 (16b)

The above parameterizations induce relative errors of no more than a few percents for both water and solute activities. The largest errors occur at low molality regions where ion activities usually exhibit a com-

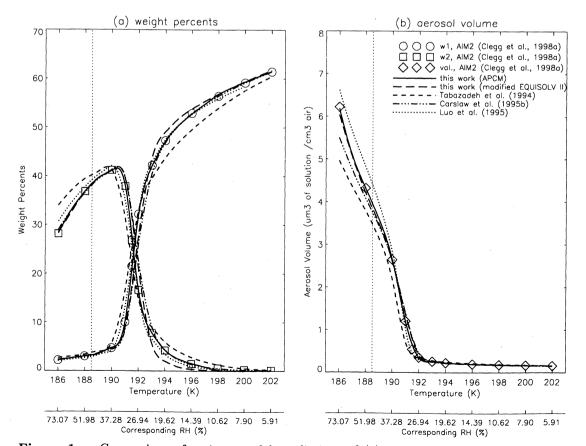


Figure 1. Comparison of various model predictions of (a) weight percents and (b) aerosol volume under background conditions in the stratosphere (see Table 7 for model conditions).

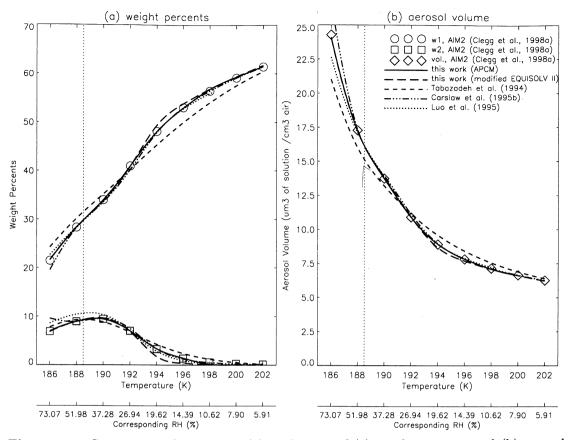


Figure 2. Comparison of various model predictions of (a) weight percents and (b) aerosol volume under volcanic conditions in the stratosphere (see Table 7 for model conditions).

plicated behavior. In addition to the improvement of binary water and solute activity coefficients, equilibrium constants of equations (11)-(13) were obtained from *Clegg et al.* [1998a] and converted into suitable molality-based units for use in EQUISOLV II. These modifications ensure perfect agreement of EQUI-SOLV II and AIM2 at a binary level where electrolyte solutions are made of  $H_2O$  plus only one electrolyte.

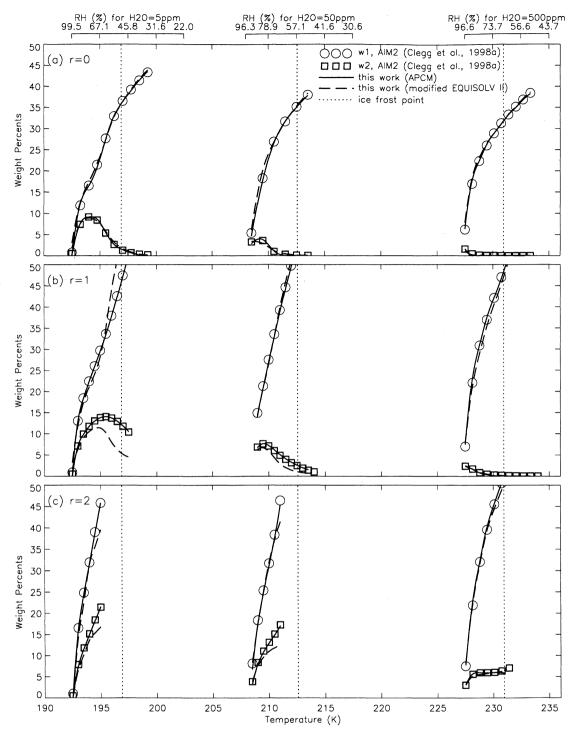
## 4. Model Intercomparison and Evaluation

#### 4.1. $H_2SO_4/HNO_3/H_2O$

For the stratospheric system of  $H_2SO_4/HNO_3/H_2O$ , the compositions obtained from APCM are compared with previous formulations [*Carslaw et al.*, 1995b; *Luo et al.*, 1995; *Tabazadeh et al.*, 1994] under background and volcanic states (see Table 7 for model conditions used). Results from model intercomparisons between six different ternary aerosol models are shown in Figures 1 and 2. For the background stratosphere, APCM with Diff correction terms included is in good agreement with AIM2 [*Carslaw et al.*, 1995a]. The predicted weight percents by APCM overlap with AIM2 almost exactly throughout the entire temperature range shown in Figure 1a, whereas previous models deviate from AIM2 either at the initial, in the middle, or at the final stage of aerosol growth (Figure 1b). Under volcanic conditions ( $H_2SO_4$  increases from 0.5 to 20ppb), all models produce nearly identical results with only slight variations (Figures 2a and 2b).

#### 4.2. $NH_3/H_2SO_4/HNO_3/H_2O$

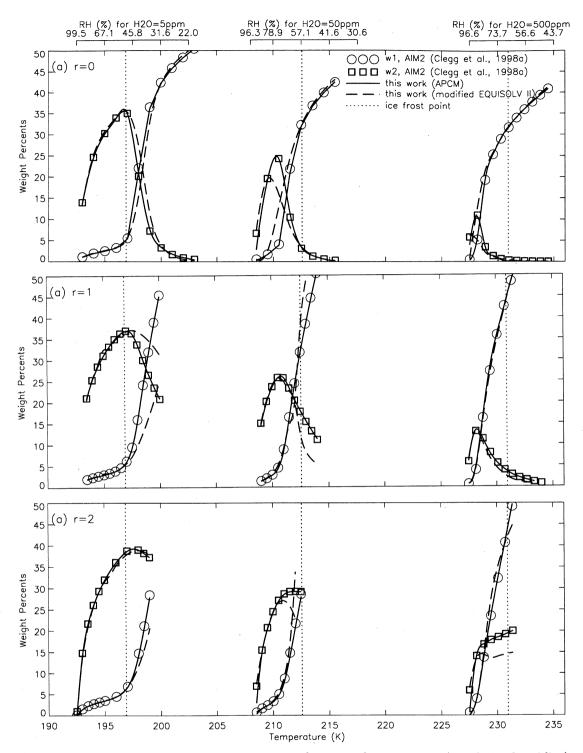
In the upper troposphere, aerosol compositions are examined for both background and polluted states (see Table 7 for model conditions). Under polluted conditions, HNO<sub>3</sub> mixing ratios are elevated well beyond background values of 100ppt [Laaksonen et al., 1997], mainly because of convective transport of polluted boundary layer air directly into the upper troposphere. Variations in the aerosol composition predicted by three models, APCM, modified EQUISOLV II (with the improved activity data), and AIM2 [Clegg et al., 1998a], are illustrated in Figures 3 and 4. The points of ice saturation are marked as vertical dotted lines in Figures 3 and 4. For all cases studied and compared, APCM yields compositions that are in close agreement with AIM2. The modified EQUISOLV II with new activity data performs better in the regions where relative humidity (RH) is high. The deviations at lower relative humidity are likely caused by the uncertainties in mixed activity coefficients calculated from simple mix-



**Figure 3.** Variation of aerosol compositions as a function of temperature (or relative humidity) under background conditions in the upper troposphere (see Table 7 for model conditions). (a) r = 0.0, (b) r = 1.0, and (c) r = 2.0.

ing rules [Bromley, 1973; Kusik and Meissner, 1978] in EQUISOLV II.

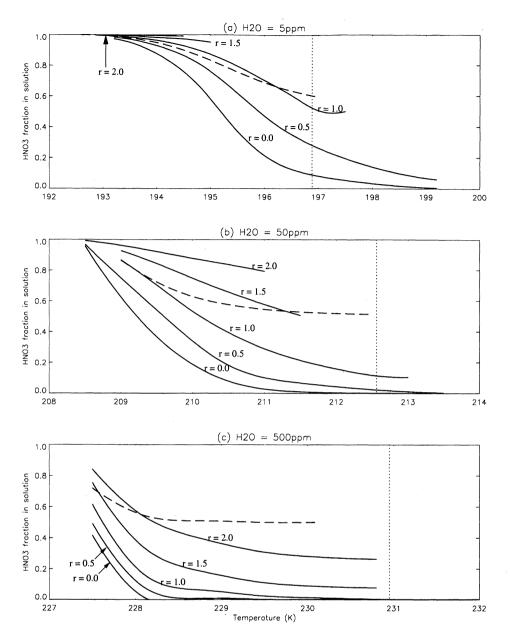
We have also compared predicted molalities of individual components,  $NH_4^+$ ,  $H_2SO_4$  (=  $HSO_4^- + SO_4^{2-}$ ), and  $NO_3^-$ , in solution (instead of weight percents of  $w_1$  and  $w_2$ ) and found that the differences between APCM and AIM2 are no more than a few percent except at the very low or high solute concentration regions. In other words, predictions by APCM are less accurate only in regions where weight percents approach the lower or upper limits (1% and 85%). Nevertheless, for very low weight percents, concentrations are too dilute to be significant, whereas for very high weight percents, relative humidities are sufficiently low that ammoniated salts, such as  $(NH_4)_2SO_4$  or  $(NH_4)_3H(SO_4)_2$ (letovicite), would most likely precipitate in solution



**Figure 4.** Variation of aerosol compositions as a function of temperature (or relative humidity) under polluted conditions in the upper troposphere (see Table 7 for model conditions). (a) r = 0.0, (b) r = 1.0, and (c) r = 2.0.

[Tabazadeh and Toon, 1998]. For example, crystallization of  $(NH_4)_2SO_4$  occurs at ~ 35% RH at room temperature [Xu et al., 1998]. At colder temperatures, crystallization of  $(NH_4)_2SO_4$  will probably occur at slightly higher RHs. In addition, for very concentrated solutions (w<sub>1</sub> + w<sub>2</sub> > 85%), aerosol compositions predicted by AIM2 are often outside the range of model validation [Clegg et al., 1998a]. For the reasons mentioned above, the limits imposed on the weight percents in the APCM are not a serious drawback. Also, as shown in Figures 3 and 4, the APCM compositions are most accurate near regions where ice reaches saturation in the atmosphere.

In addition to aerosol composition the APCM can be used to calculate the extent of  $HNO_3$  uptake by up-



**Figure 5.** Uptake of  $HNO_3$  in sulfate-based aerosols as a function of sulfate neutralization under background conditions in the upper troposphere (see Table 7 for model conditions). (a)  $H_2O = 5ppm$ , (b)  $H_2O = 50ppm$ , and (c)  $H_2O = 500ppm$ . (Dashed lines represent the uptake by pure NH<sub>4</sub>NO<sub>3</sub> aerosols (i.e.,  $H_2SO_4 = 0$ ))

per tropospheric aerosols. Aircraft field experiments have shown that ammoniated particles are abundant in the upper troposphere [*Talbot et al.*, 1996, 1998; *Tabazadeh et al.*, 1998]. Here we examine how HNO<sub>3</sub> uptake may be affected by the presence of ammoniated aerosols in the upper troposphere. For simulations the H<sub>2</sub>SO<sub>4</sub> and HNO<sub>3</sub> mixing ratios are fixed at background levels of 100 ppt each. NH<sub>3</sub> concentrations in the APCM are varied by increasing (or decreasing) r to account for changes in solution neutrality (or pH), ranging from pure H<sub>2</sub>SO<sub>4</sub> solution droplets (r = 0.0) to fully ammoniated systems consisting of only (NH<sub>4</sub>)<sub>2</sub>SO<sub>4</sub> (r = 2.0) (see Table 7 for model conditions used). In Figure 5, the uptake of HNO<sub>3</sub>, expressed as a fraction (i.e., ratio of concentration in the liquid phase to the total initial concentration), is shown for three different assumed water vapor pressure profiles. As expected, partitioning of HNO<sub>3</sub> in sulfate-based aerosols depends strongly on solution neutrality. In general, greater uptake occurs at lower temperatures and higher relative humidities. For example, at temperatures lower than  $210^{\circ}$ K (which corresponds to a relative humidity of 78% for P<sub>total</sub> = 200mb and H<sub>2</sub>O = 50ppm), a significant fraction (> 90%) of HNO<sub>3</sub> resides in fully neutralized ammoniated solutions, compared to only < 20% in the pure sulfate system (Figure 5b). It has been shown that high levels of  $HNO_3$  in solution may cause precipitation of ammoniated and/or nitrated salts, which could change the mode of ice formation from homogeneous to heterogeneous nucleation [*Tabazadeh and Toon*, 1998].

#### 5. Conclusions

The ion interaction model of Clegg et al. [1998a] for the system of  $NH_3/H_2SO_4/HNO_3/H_2O$  has been parameterized into a compact model well suited for incorporation into large-scale atmospheric models. The aerosol physical chemistry model (APCM) reproduces the AIM2 results of Clegg et al. [1998a] for a wide range of conditions in the upper troposphere and lower stratosphere. Model intercomparisons show that for the ternary system of H<sub>2</sub>SO<sub>4</sub>/HNO<sub>3</sub>/H<sub>2</sub>O solution compositions obtained from APCM are in better agreement with those obtained from AIM2 than previous formulations. For the quaternary system of NH<sub>3</sub>/H<sub>2</sub>SO<sub>4</sub>/HNO<sub>3</sub> /H<sub>2</sub>O, APCM results are also in good agreement with AIM2 predictions, particularly near the regions of ice saturation where the influence of ammoniated particles on the ice nucleation process is of interest. Extension of APCM to include other features such as calculations of deliquescence relative humidity and precipitation of solids in solution are under way.

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#### References

- Adams, P. J., J. H. Seinfeld, and D. M. Koch, Global concentrations of tropospheric sulfate, nitrate, and ammonium aerosol simulated in a general circulation model, J. Geophys. Res., 104, 13,791, 1999.
- Bassett, M., and J. H. Seinfeld, Atmospheric equilibrium model of sulfate and nitrate aerosols, Atmos. Environ., 17, 2237, 1983.
- Bromley, L. A., Thermodynamic properties of strong electrolytes in aqueous solutions, AIChE J., 19, 313, 1973.
- Carslaw, K. S., B. P. Luo, S. L. Clegg, T. Peter, P. Brimblecombe, and P. Crutzen, Stratospheric aerosol growth and HNO<sub>3</sub> gas phase depletion from coupled HNO<sub>3</sub> and water uptake by liquid particles, *Geophys. Res. Lett.*, 21, 2479, 1994.
- Carslaw, K. S., S. L. Clegg, and P. Brimblecombe, A thermodynamic model of the system  $HCl-HNO_3-H_2SO_4-H_2O$ , including solubilities of HBr, from 328K to < 200K. J. Phys. Chem., 99, 11,557, 1995a.
- Carslaw, K. S., B. P. Luo, and T. Peter, An analytic expression for the composition of aqueous HNO<sub>3</sub>-H<sub>2</sub>SO<sub>4</sub> strato-

spheric aerosols including gas phase removal of HNO<sub>3</sub>, Geophys. Res. Lett., 22, 1877, 1995b.

- Clegg., S. L., and P. Brimblecombe, Equilibrium partial pressures and mean activity and osmotic coefficients of 0-100% nitric acid as a function of temperature, J. Phys. Chem., 94, 5369, 1990.
- Clegg., S. L., and P. Brimblecombe, A generalized multicomponent thermodynamic model applied to the (NH<sub>4</sub>)<sub>2</sub>SO<sub>4</sub>-H<sub>2</sub>SO<sub>4</sub>-H<sub>2</sub>O system to high supersaturation and low relative humidity at 298.15K, J. Aerosol Sci., 26, 19, 1995a.
- Clegg., S. L., and P. Brimblecombe, Application of a multicomponent thermodynamic model to activities and thermal properties of 0-40 mol kg<sup>-1</sup> aqueous sulfuric acid from < 200 to 328K, J. Chem. Eng. Data, 40, 43, 1995b.
- Clegg, S. L., P. Brimblecombe, and A. S. Wexler, A thermodynamic model of the system  $H^+ - NH_4^+ - O_4^{-2} - NO_3^- - H_2O$  at tropospheric temperatures, J. Phys. Chem. A, 102, 2137, 1998a.
- Clegg, S. L., P. Brimblecombe, and A. S. Wexler, A thermodynamic model of the system  $H^+ - NH_4^+ - Na^+ - SO_4^{-2} - NO_3^- - Cl^- - H_2O$  at 298.15K, J. Phys. Chem. A, 102, 2155, 1998b.
- Jacobson, M. Z., Fundamentals of Atmospheric Modeling, Cambridge Univ. Press, New York, 1999a.
- Jacobson, M. Z., Studying the effects of calcium and magnesium on size-distributed nitrate and ammonium with EQUISOLV II, Atmos. Environ., 33, 3635, 1999b.
- Jacobson, M. Z., A physically-based treatment of elemental carbon optics: Implications for global direct forcing of aerosols, *Geophys. Res. Lett.*, 27, 217, 2000.
- Jacobson, M. Z., A. Tabazadeh, and R. P. Turco, Simulating equilibrium within aerosols and nonequilibrium between gases and aerosols, J. Geophys. Res., 101, 9079, 1996.
- Kärcher, B., and S. Solomon, On the composition and optical extinction of particles in the tropopause region, J. Geophys. Res., 104, 27,441, 1999.
- Kim, Y. P., and J. H. Seinfeld, Atmospheric gas-aerosol equilibrium, III, Thermodynamics of crustal elements Ca<sup>2+</sup>, K<sup>+</sup>, and Mg<sup>2+</sup>, Aerosol Sci. Technol., 22, 93, 1995.
- Kusik, C. L., and H. P. Meissner, Electrolyte activity coefficients in inorganic processing, AIChE Symp. Ser., 74, 14, 1978.
- Laaksonen, A., J. Hienola, and M. Kulmala, Supercooled cirrus cloud formation modified by nitric acid pollution of the upper troposphere, *Geophys. Res. Lett.*, 24, 3009, 1997.
- Luo, B., K. S. Carslaw, T. Peter, and S. L. Clegg, Vapor pressures of H<sub>2</sub>SO<sub>4</sub>/HNO<sub>3</sub>/HCl/HBr/H<sub>2</sub>O solutions to low stratospheric temperatures, *Geophys. Res. Lett.*, 22, 247, 1995.
- Pilinis, C., and J. H. Seinfeld, Continued development of a general equilibrium model for inorganic multicomponent atmospheric aerosols, Atmos. Environ., 32, 2453, 1987.
- Pitzer, K. S., Activity Coefficients in Electrolyte Solutions, 2nd ed., CRC Press, Boca Raton, Fla., 1991.
- Saxena, P., C. Seigneur, A. B. Hudischewskyj, and J. H. Seinfeld, A comparative study of equilibrium approaches to the chemical characterizations of secondary aerosols, *Atmos. Environ.*, 20, 1471, 1986.
- Solomon, S., Stratospheric ozone depletion: A review of concepts and history, Rev. Geophys., 37, 275, 1999.
- Stelson, A. W., and J. H. Seinfeld, Relative humidity and temperature dependence of the ammonium nitrate dissociation constant, Atmos. Environ., 26, 983, 1981.
- Stokes, R. H., and R. A. Robinson, Interactions in aqueous nonelectrolyte solutions, I, solute-solvent equilibria, J. Phys. Chem., 70, 2126, 1966.

- Tabazadeh, A., and O. B. Toon, The role of ammoniated aerosols in cirrus cloud nucleation, *Geophys. Res. Lett.*, 25, 1379, 1998.
- Tabazadeh, A., R. P. Turco, and M. Z. Jacobson, A model for studying the composition and chemical effects of stratospheric aerosols, J. Geophys. Res., 99, 12,897, 1994.
- Tabazadeh, A., O. B. Toon, S. Clegg, and P. Hamill, A new parameterization of H<sub>2</sub>SO<sub>4</sub>/H<sub>2</sub>O aerosol composition: Atmospheric implications, *Geophys. Res. Lett.*, 24, 1931, 1997a.
- Tabazadeh, A., O. B. Toon, and E. Jensen, Formation and implications of ice particle nucleation in the stratosphere, *Geophys. Res. Lett.*, 24, 2007, 1997b.
- Tabazadeh, A., M. Z. Jacobson, H. B. Singh, O. B. Toon, J. S. Lin, B. Chatfield, A. N. Thakur, R. W. Talbot, and J. E. Dibb, Nitric acid scavenging by mineral and biomass aerosols, *Geophys. Res. Lett.*, 25, 4185, 1998.
- Talbot, R. W. et al., Chemical characteristics of continental outflow from Asia to the troposphere over the western Pacific Ocean during September-October 1991: Results from PEM-West A, J. Geophys. Res., 101, 1713, 1996.

- Talbot, R. W., J. E. Dibb, and M. B. Loomis, Influence of vertical transport on free tropospheric aerosols over the central USA in springtime, *Geophys. Res. Lett.*, 25, 1367, 1998.
- Weisenstein, D. K., Y. K. Glenn, M. K. W. Ko, N. Sze, J. M. Rodriguez, and C. J. Scott, A two-dimensional model of sulfur species and aerosols, J. Geophys. Res., 102, 13,019, 1997.
- Wexler, A. S., and J. H. Seinfeld, Second-generation inorganic aerosol model, Atmos. Environ., Part A, 25, 2731, 1991.
- Xu, J., D. Imre, R. McGraw, and I. Tang, Ammonium sulfate: Equilibrium and metastability phase diagrams from  $40^{\circ}$ C to  $-50^{\circ}$ C, J. Phys. Chem. B, 102, 7462, 1998.

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