

ANALYSIS OF A "SLUG TEST" OR DRILL STEM TEST FOR LINEAR FLOW

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INTRODUCTION

The objective of this work is to provide a means for analyzing pressure transients from drill stem tests (DSTs) in fractured wells dominated by linear flow in the formation. The consequent partial differential equations have been solved by numerical inversion of the Laplace transform. A DST consists of isolating an open stretch of borehole with packers and connecting it instantaneously with the atmosphere by means of a drill-string. In water-dominated reservoirs, at the end of the test a level usually stabilizes in the string and the well does not produce spontaneously. The method can be applied only to the "analysis of DSTs in which the fluid influx into the drill pipe tends to kill the flow, giving only a partial DST recovery."¹ The analysis is performed by type-curve matching, and leads to the determination of the initial reservoir pressure, flow characteristics of the formation, and wellbore damage. These parameters allow the prediction of early-time pressure drawdown due to a constant flowrate.

Methods already exist in the literature¹ for analyzing these pressure transients in the presence of radial flow in the formation. In the fractured reservoirs found in some geothermal² fields, linear flow models seem more appropriate than radial models.

DESCRIPTION OF THE MODEL AND RESULTS

Figure 1 presents the assumed physical model. It consists of:

a) A homogeneous one-dimensional rock formation in which the pressure propagates according to the diffusion equation:

$$\frac{\partial^2 v(x,t)}{\partial x^2} = \frac{1}{\eta} \frac{\partial v(x,t)}{\partial t} \quad (1)$$

where: $v(x,t)$ = pressure drawdown in the formation
 x = distance from well
 t = time
 η = formation diffusivity, $k/\phi\mu c$

Editor's Note: this paper was presented at the Fourth Workshop, but inadvertently omitted from the Proceedings.

Equation 1 describes fluid flow in a porous medium according to Darcy's law and the continuity equation, and assuming the fluid has a small constant compressibility.³

b) A concentrated flow resistance, R, that produces a pressure drawdown proportional to the flowrate through it. This could, for example, be caused by a layer of mud on the walls of the bore. The following equation therefore holds:

$$Rq(t) = u(t) - v(o,t) \quad (2)$$

where $q(t)$ = instantaneous volumetric flowrate
 $u(t)$ = pressure drawdown at bottomhole

c) The fluid produced by the formation is stored in the drill-string, causing a back-pressure on the formation proportional to the fluid stored. Therefore, we have:

$$-C \frac{du(t)}{dt} = q(t) \quad (3)$$

where C is a constant:

$$C = \frac{\text{internal cross-section of drill pipe}}{\text{specific weight of fluid}} = \frac{A}{\gamma}$$

The partial differential equation (Eq. 1), along with the boundary conditions (Eqs. 2 and 3), Eq. A6 (in the Appendix), and with the initial conditions:

$$v(x,0) = 0$$

$$u(o) = V$$

$$x > 0$$

was solved by means of the Laplace transform, as in Carslaw and Jaeger.⁴ The solution is:

$$\frac{\bar{u}(p)}{V} = \frac{\sigma\sqrt{p} + 1}{\sqrt{p}(\sigma\sqrt{p} + 1)} \quad (4)$$

where $\bar{u}(p)$ is the Laplace transform of $u(t)$ and p its parameter, while σ is a dimensionless constant:

$$\sigma = \beta CR$$

$$\text{where } \beta = \left(\frac{Sk}{\mu C} \right)^2 \frac{1}{\eta}$$

See the Appendix for further details.

Unfortunately, the analytical inversion of Eq. 4 is not straightforward, so the numerical method given in Eq. 5 is used. The results are given in Tables 1 and 2 for different values of σ , and summarized in the type-curves 1, 2, and 3. The computer code is also included for the numerical inversion of Eq. 4.

USING THE TYPE-CURVES

a) If the initial reservoir pressure is unknown, type-curves 1 and 2 should be used. These show:

$$\frac{V-u(t)}{V} = 1 - \frac{u(t)}{V}$$

as a function of $t_D = \beta t$ on a log-log graph. The pressure recovery, $V-u(t)$, should be plotted against time, t , on log-log tracing paper of the same scale. The field data graph should be shifted without rotation over the type-curves until the best match is obtained. V can be estimated; hence, initial reservoir pressure from the vertical match.

From the horizontal match, we can estimate the value of the constant:

$$\beta = \frac{t_D}{t} = \left(\frac{Sk}{\mu C} \right)^2 \cdot \frac{1}{\eta}$$

From the selected type-curve, we obtain:

$$\sigma = \beta CR$$

3) If the initial reservoir pressure is known or was estimated as described in a), type-curve 3 should be used, since it permits a better evaluation of the parameters. In this case, semilog tracing paper is used, plotting $u(t)/V$ versus time, t . Once again, the scales should be the same as those of the type-curve.

In this case, we can obtain the match only by shifting the data horizontally over the type-curve, obtaining:

$$\beta = \frac{t_D}{t} = \left(\frac{Sk}{\mu C} \right)^2 \cdot \frac{1}{\eta}$$

and:

$$\sigma = \beta R$$

In both cases a) and b), having evaluated σ and β and knowing $C = A/\gamma$, we obtain R and the group $(Sk/\mu)^2 \cdot 1/\eta$, which controls linear flow.

From the above parameters, it is possible to forecast the early-time pressure transient caused by a constant flowrate. In fact, from Eq. 6:

$$v(o, t) = \frac{2q\sqrt{\eta}}{\sqrt{\pi}Sk\mu} \cdot \sqrt{t} = \frac{2}{\sqrt{\pi}} \cdot \frac{q}{C\sqrt{\beta}} \cdot \sqrt{t} \quad (5)$$

Hence:

$$u(t) = Rq + \frac{2}{\sqrt{\pi}} \cdot \frac{q}{c\sqrt{\beta}} \cdot \sqrt{t} = q\beta c \left(\sigma + \frac{2}{\sqrt{\pi}} \sqrt{\beta t} \right) \quad (6)$$

All the parameters in Eq. 6 are either known or determined by this analysis. Where R is too high, i.e., where σ is large in comparison to $2/\sqrt{\pi} \sqrt{\beta t}$, stimulation could prove useful.

ACKNOWLEDGEMENTS

Thanks are extended to P. G. Atkinson and H. J. Ramey, Jr., for providing the bibliography and subroutine for the numerical inversion of the Laplace transform.

NOMENCLATURE

- A = internal cross-section of drill pipe, m^2
- c = effective compressibility, Pa^{-1}
- $C = \frac{A}{\gamma}$ = fluid storage constant in drill string, $m^4 kg^{-1} s^2$
- D(p) = function independent of x, Pa
- E(p) = function independent of x, Pa
- k = formation permeability, m^2
- p = Laplace transform parameter
- q(t) = volumetric flowrate from formation to drill pipe, $m^3 s^{-1}$
- R = concentrated flow resistance, $m^{-4} kg s^{-1}$
- S = fracture surface crossed by flowrate q. If fluid flows from opposite directions, this surface should be doubled, m^2
- t = time from when bottomhole valve is opened, seconds
- $t_D = \beta t$ = dimensionless time
- u(t) = pressure drawdown in wellbore, Pa
- $\bar{u}(p)$ = Laplace transform of u(t), Pa
- v(x,t) = pressure drawdown in formation, Pa
- $\bar{v}(x,p)$ = Laplace transform of v(x,t), Pa
- V=u(o) = maximum pressure drawdown in wellbore at beginning of test, Pa
- x = distance from fracture (see Fig. 1), m
- $\alpha = \frac{\beta p}{\eta} = \frac{\sigma}{\epsilon} \sqrt{p}$ = a constant with respect to x and t, m^{-1}
- $\beta = \frac{Sk}{\mu C} \cdot \frac{1}{\eta} = \frac{t_D}{t}$ = a constant, s^{-1}
- γ = specific weight of fluid produced, $m^{-2} kg s^{-2}$

$\epsilon = \frac{SkR}{\mu}$ = formation thickness equivalent to concentrated resistance, m
 μ = fluid viscosity, Pa(s)
 σ = βCR = a constant
 $\eta = \frac{k}{\phi \mu C}$ = formation diffusivity, $m^2 s^{-1}$
 ϕ = porosity

REFERENCES

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APPENDIX

SOLUTION OF THE PARTIAL DIFFERENTIAL EQUATION

Introducing dimensionless time, $t_D = \beta t$, Eq. 1 becomes:

$$\frac{\partial^2 v(x, t_D)}{\partial x^2} = \frac{\beta}{\eta} \frac{\partial v(x, t_D)}{\partial t_D} \quad (A-1)$$

Laplace transforming and remembering that $v(x, 0) = 0$, we have:

$$\frac{\partial^2 \bar{v}(x, p)}{\partial x^2} = \frac{\beta p}{\eta} \cdot \bar{v}(x, p) \quad (A-2)$$

where $\bar{v}(x, p)$ is the Laplace transform of $v(x, t_D)$ and p its parameter.

Combining Eqs. 2 and 3, we obtain the first boundary condition at $x = 0$:

$$C \frac{du(t)}{dt} + \frac{u(t) - v(0, t)}{R} = 0 \quad (A-3)$$

Introducing dimensionless time, t_D , and the constant $\sigma = \beta CR$ into Eq. A-3, we get:

$$\sigma \frac{du(t_D)}{dt_D} + u(t_D) - v(0, t_D) = 0 \quad (A-4)$$

Laplace transforming and defining the initial pressure disturbance, $V = u(0)$, we obtain:

$$(\sigma p + 1) \cdot \bar{u}(p) - \bar{v}(0, p) = \sigma V \quad (A-5)$$

where $\bar{u}(p)$ is the Laplace transform of $u(t_D)$.

The second boundary condition at $x = 0$ derives from Darcy's law:

$$q(t) = \frac{Sk}{\mu} \left(\frac{\partial v(x, t)}{\partial x} \right)_{x=0} \quad (A-6)$$

which, combined with Eq. 2, becomes:

$$\frac{u(t) - v(0, t)}{R} + \frac{Sk}{\mu} \left(\frac{\partial v(x, t)}{\partial x} \right)_{x=0} = 0 \quad (A-7)$$

Defining $\epsilon = \frac{SkR}{\mu}$, introducing dimensionless time, t_D , and Laplace transforming Eq. A-7, we have:

$$\bar{u}(p) - \bar{v}(0,p) + \epsilon \left(\frac{\partial v(x,p)}{\partial x} \right)_{x=0} = 0 \quad (A-8)$$

Now we must solve Eq. A-2 with Boundary conditions (Eqs. A-5 and A-8). The general solution of Eq. A-2 is:

$$\bar{v}(x,p) = D(p) e^{-\alpha x} + E(p) e^{+\alpha x} \quad (A-9)$$

with $\alpha = \sqrt{\frac{\beta p}{\eta}} = \frac{\sigma}{\epsilon} \sqrt{p}$ and $D(p)$ and $E(p)$ two functions non x -dependent.

Since $v(x,t)$, and hence $v(x,p)$ are limited for $x \rightarrow \infty$, $E(p)$ must be nil and the solution we are looking for is:

$$\bar{v}(x,p) = D(p) e^{-\alpha x} \quad (A-10)$$

Substituting Eq. A-10 into Eqs. A-5 and A-8, we get the expression for $D(p)$ and $\bar{u}(p)$:

$$D(p) = \frac{V}{\sqrt{p}(\sigma\sqrt{p}+1)} \quad (A-11)$$

$$\bar{u}(p) = D(p) \cdot (\sigma\sqrt{p} + 1) \quad (A-12)$$

Hence:

$$\frac{\bar{u}(p)}{V} = \frac{\sigma\sqrt{p} + 1}{\sqrt{p}(\sigma\sqrt{p}+1)} \quad (A-13)$$

TABLE 1

-20-

t_D	$\frac{u}{V}$		
	$\sigma = 0$	$\sigma = 0.01$	$\sigma = 0.1$
0.158490-05	0.998580+00	0.999860+00	
0.251190-05	0.998210+00	0.999770+00	
0.398110-05	0.997750+00	0.999660+00	
0.630960-05	0.997170+00	0.999470+00	
0.100000-04	0.996440+00	0.999200+00	
0.158490-04	0.995530+00	0.998790+00	
0.251190-04	0.994370+00	0.998200+00	
0.398110-04	0.992920+00	0.997340+00	
0.630960-04	0.991100+00	0.996140+00	
0.100000-03	0.988820+00	0.994460+00	
0.158490-03	0.985950+00	0.992180+00	0.998560+00
0.251190-03	0.982370+00	0.989110+00	0.997760+00
0.398110-03	0.977880+00	0.985080+00	0.996550+00
0.630960-03	0.972280+00	0.979840+00	0.994730+00
0.100000-02	0.965290+00	0.973120+00	0.992000+00
0.158490-02	0.956620+00	0.964620+00	0.987970+00
0.251190-02	0.945870+00	0.953940+00	0.982120+00
0.398110-02	0.932600+00	0.940640+00	0.973780+00
0.630960-02	0.916320+00	0.924230+00	0.962130+00
0.100000-01	0.896460+00	0.904140+00	0.946230+00
0.158490-01	0.872410+00	0.879760+00	0.925070+00
0.251190-01	0.843580+00	0.850500+00	0.897660+00
0.398110-01	0.809400+00	0.815800+00	0.863180+00
0.630960-01	0.769460+00	0.775250+00	0.821080+00
0.100000+00	0.723580+00	0.728670+00	0.771240+00
0.158490+00	0.671910+00	0.676260+00	0.714130+00
0.251190+00	0.615080+00	0.618670+00	0.650830+00
0.398110+00	0.554250+00	0.557090+00	0.583060+00
0.630960+00	0.491070+00	0.493220+00	0.513110+00
0.100000+01	0.427580+00	0.429130+00	0.443530+00
0.158490+01	0.365970+00	0.367030+00	0.376870+00
0.251190+01	0.308220+00	0.308910+00	0.315260+00
0.398110+01	0.255900+00	0.256320+00	0.260210+00
0.630960+01	0.209930+00	0.210170+00	0.212440+00
0.100000+02	0.170580+00	0.170720+00	0.171990+00
0.158490+02	0.137620+00	0.137690+00	0.138380+00
0.251190+02	0.110450+00	0.110490+00	0.110860+00
0.398110+02	0.883350-01	0.883560-01	0.885450-01
0.630960+02	0.704770-01	0.704880-01	0.705850-01
0.100000+03	0.561410-01	0.561470-01	0.561960-01
0.158490+03	0.446750-01	0.446780-01	0.447030-01
0.251190+03	0.355280-01	0.355290-01	0.355410-01
0.398110+03	0.282410-01	0.282420-01	0.282480-01
0.630960+03	0.224430-01	0.224430-01	0.224470-01
0.100000+04	0.178320-01	0.178330-01	0.178340-01
0.158490+04	0.141670-01	0.141670-01	0.141680-01
0.251190+04	0.112550-01	0.112550-01	0.112550-01
0.398110+04	0.894070-02	0.894080-02	0.894090-02
0.630960+04	0.710220-02	0.710210-02	0.710230-02
0.100000+05	0.564160-02	0.564160-02	0.564170-02

TABLE 2

-21-

t_D	$\frac{u}{v}$		
	$\sigma = 1$	$\sigma = 10$	$\sigma = 100$
0.158490-01	0.985640+00	0.998430+00	0.999840+00
0.251190-01	0.977850+00	0.997520+00	0.999750+00
0.398110-01	0.966080+00	0.996090+00	0.999600+00
0.630960-01	0.948570+00	0.993830+00	0.999370+00
0.100000+00	0.923000+00	0.990280+00	0.999000+00
0.158490+00	0.886610+00	0.984730+00	0.998420+00
0.251190+00	0.836510+00	0.976090+00	0.997500+00
0.398110+00	0.770520+00	0.962730+00	0.996050+00
0.630960+00	0.688380+00	0.942260+00	0.993750+00
0.100000+01	0.593240+00	0.911360+00	0.990120+00
0.158490+01	0.492170+00	0.865660+00	0.984420+00
0.251190+01	0.394740+00	0.800160+00	0.975480+00
0.398110+01	0.309430+00	0.710430+00	0.961540+00
0.630960+01	0.240360+00	0.595440+00	0.939970+00
0.100000+02	0.186920+00	0.461640+00	0.906990+00
0.158490+02	0.146110+00	0.325810+00	0.857510+00
0.251190+02	0.114800+00	0.210910+00	0.785390+00
0.398110+02	0.905380-01	0.132560+00	0.684870+00
0.630960+02	0.715890-01	0.880250-01	0.554050+00
0.100000+03	0.567010-01	0.633900-01	0.400780+00
0.158490+03	0.449560-01	0.479350-01	0.247650+00
0.251190+03	0.356690-01	0.370700-01	0.126630+00
0.398110+03	0.283120-01	0.289890-01	0.575180-01
0.630960+03	0.224790-01	0.228100-01	0.300730-01
0.100000+04	0.178500-01	0.180140-01	0.203040-01
0.158490+04	0.141760-01	0.142580-01	0.152210-01
0.251190+04	0.112600-01	0.113000-01	0.117500-01
0.398110+04	0.894300-02	0.896320-02	0.917950-02
0.630960+04	0.710330-02	0.711340-02	0.721880-02
0.100000+05	0.564230-02	0.564730-02	0.569930-02

COMPUTER CODE

C QUESTO PROGRAMMA INVERTE LA TRASFORMATATA DI LAPLACE

```

      IMPLICIT REAL*8(A-H,O-Z)
      COMMON G(50),V(50),H(25),M
      N=18
      T=0.01
      WRITE(6,2)N
2     FORMAT(7X,'T',15X,'FA',7X,'N=',I2)
10    T=T*10.** (1./5.)
      CALL LINV(T,FA,N)
      WRITE(6,1)T,FA
1     FORMAT(2E15.5)
      IF(T.LT.10000.) GO TO 10
      STOP
      END

```

C

FUNCTION P(ARG)

C QUESTA E' LA TRASFORMATATA DI LAPLACE CHE DEVE ESSERE

C DALLA SUBROUTINE LINV

```

      IMPLICIT REAL*8(A-H,O-Z)
      SIGMA=100.
      RQ=DSQRT(ARG)
      P=(SIGMA*RQ+1.)/(RQ*(SIGMA*ARG+RQ+1.))
      RETURN
      END

```

C

```

SUBROUTINE LINV(T,FA,N)
IMPLICIT REAL*8 (A-H,O-Z)
COMMON G(50), V(50), H(25), M
DLQGTW = .6931471805599453
IF(M.EQ.N)GO TO 100

```

LAP00010
LAP00020
LAP00030
LAP00040
LAP00050
LAP00060
LAP00070
LAP00080
LAP00090
LAP00100
LAP00110
LAP00120
LAP00130
LAP00140
LAP00150
LAP00160
LAP00170
LAP00180
LAP00190
LAP00200
LAP00210
LAP00220
LAP00230
LAP00240
LAP00250
LAP00260
LAP00270
LAP00280
LAP00290
LAP00300

```

C   CALCULATE V-ARRAY
      M=N
      G(1)=1.
      NH=N/2
      DO 5 I=2,N
5    G(I)=G(I-1)*I
      H(1)=2./G(NH-1)
      DO 10 I=2,NH
      FI=I
      IF (I.EQ.NH) GO TO 8
      H(I)=FI**NH*G(2*I)/(G(NH-I)*G(I)*G(I-1))
      GO TO 10
8    H(I)=FI**NH*G(2*I)/(G(I)*G(I-1))
10   CONTINUE
      SN=2*(NH-NH/2*2)-1
      DO 50 I=1,N
      V(I)=0.
      K1=(I+1)/2
      K2=I
      IF (K2.GT.NH) K2=NH
      DO 40 K=K1,K2
      IF (2*K-I.EQ.0) GO TO 37
      IF (I.EQ.K) GO TO 38
      V(I)=V(I)+H(K)/(G(I-K)*G(2*K-I))
      GO TO 40
37   V(I)=V(I)+H(K)/(G(I-K))
      GO TO 40
38   V(I)=V(I)+H(K)/G(2*K-I)
40   CONTINUE
      V(I)=SN*V(I)
      SN=-SN
50   CONTINUE
100  FA=0.
      A = DLOGTW/T
      DO 110 I=1,N
      ARG=I*A
110  FA=FA+V(I)*P(ARG)
      FA=A*FA
      RETURN
      END

```

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LAP00310
LAP00320
LAP00330
LAP00340
LAP00350
LAP00360
LAP00370
LAP00380
LAP00390
LAP00400
LAP00410
LAP00420
LAP00430
LAP00440
LAP00450
LAP00460
LAP00470
LAP00480
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LAP00660
LAP00670
LAP00680
LAP00690
LAP00700

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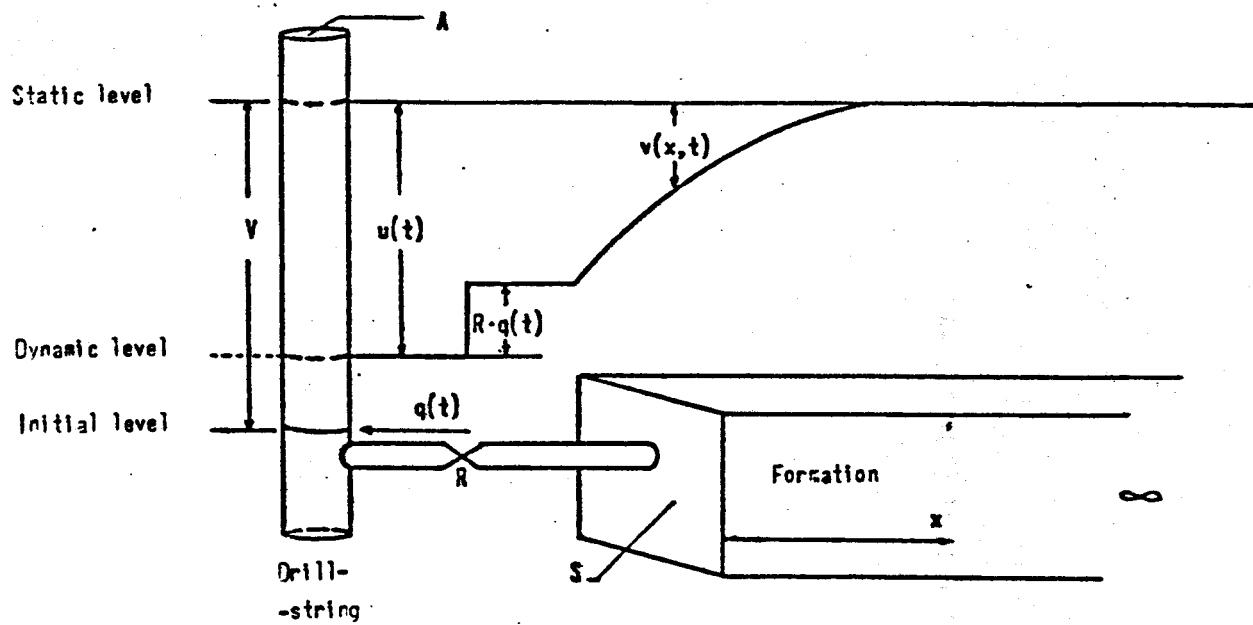
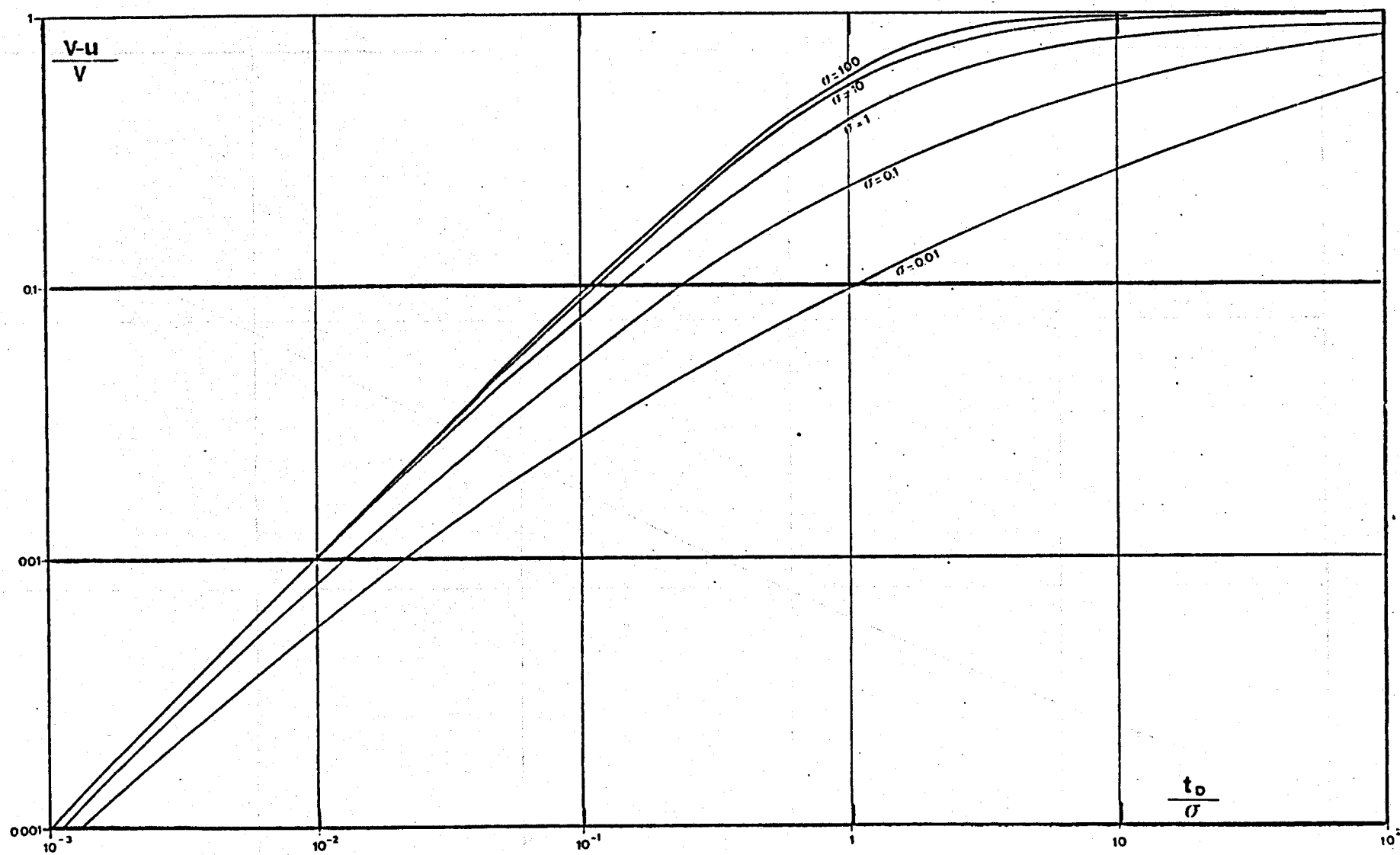
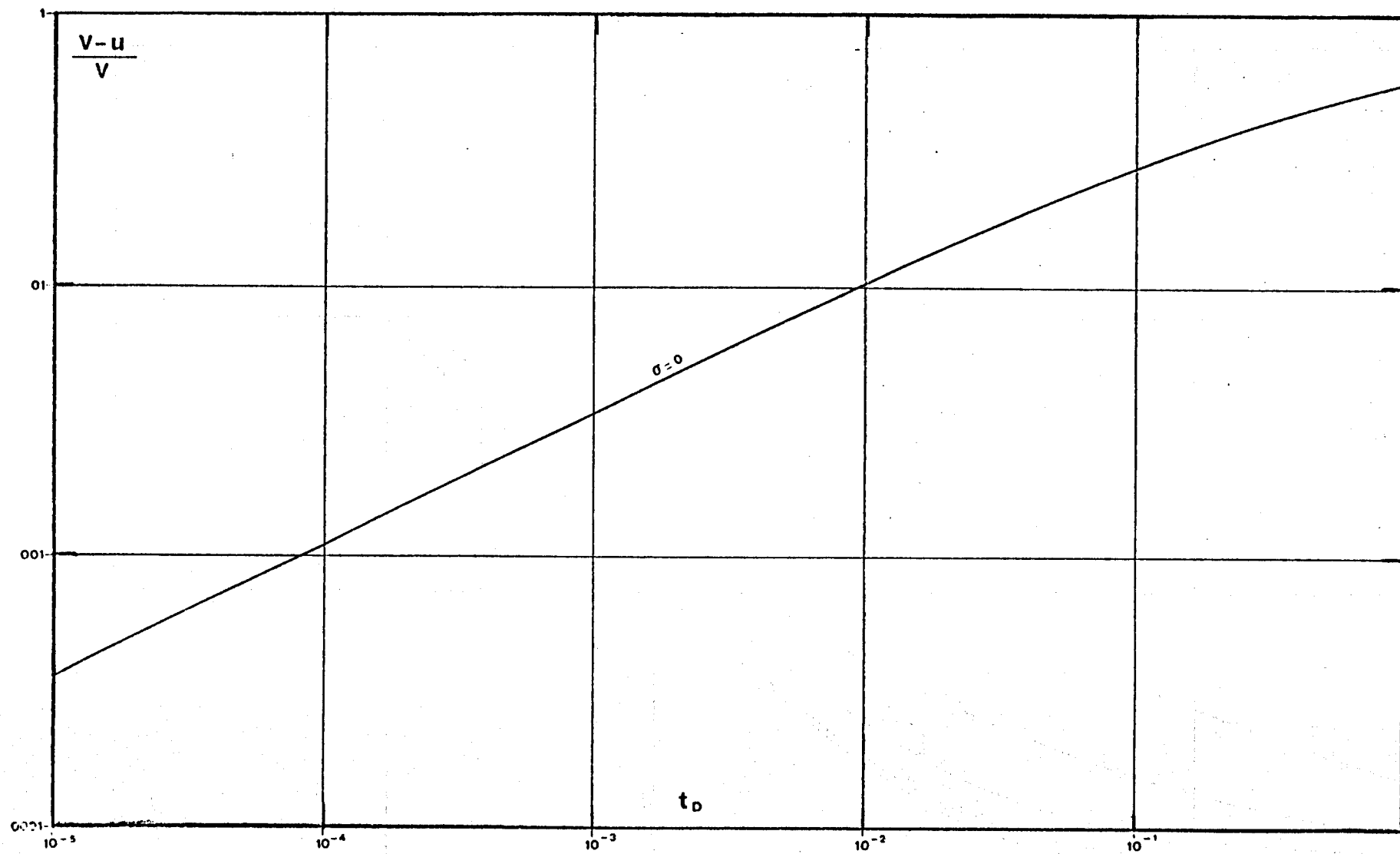


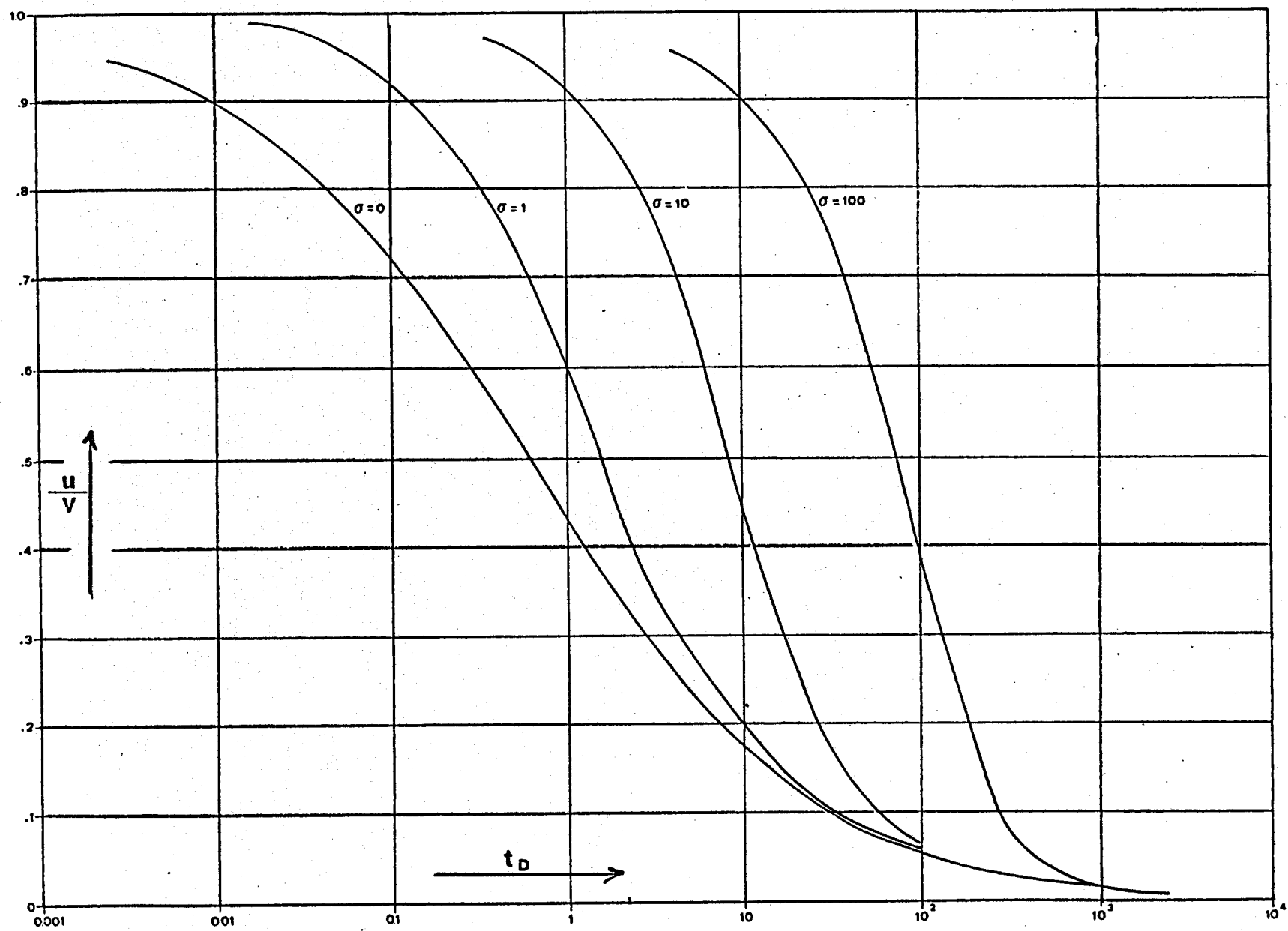
Fig. 1



Curva tipo n. 1



Curva tipo n. 2



Curva tipo n. 3